

- Outline :
- energy scales : $k_B T$ is our "ruler"
 - molecular interactions
 - thermal forces

- Readings :
- Mahadevan ch. 3
 - Daune ch. 2
 - Dill & Bromberg ch. 6, 8, 10

▷ Length scales in biology $10^{-9} \rightarrow 10^1$ m

central dogma : DNA \rightarrow RNA \rightarrow proteins \rightarrow organelles \rightarrow cells \rightarrow tissues \rightarrow organs

❖ Central problem : characterize mechanical interactions from nm to m (chemical)

- Molecular mechanics :
- conformation of macromolecules (DNA, proteins ...)
 - apply forces \Rightarrow change conformations
 - things are "noisy" at the molecular level
 - mechanochemistry

◻ Energy scales in biology

• Thermal energy : $E_{th} \sim k_B T$

\downarrow temperature
 Boltzmann constant $k_B = 1.38 \cdot 10^{-23} \text{ J.K}^{-1}$

$k_B T \approx 4 \cdot 10^{-21} \text{ J}$
 $RT = N k_B T = 25 \text{ kJ.mol}^{-1}$

at room temperature 298 K
in terms of mols

• Comparison of energies ($\times 10^{-21} \text{ J}$)

E_{th}	kT	4.1	
ATP hydrolysis	ΔG	~ 100	
electron transport	eV	~ 29	(for 180 mV)
covalent bond	C-H	~ 660	
non-covalent interactions		$\sim 0.1 - 60$	

↳ hydrogen bonds, hydrophobic interaction, screened Coulombic interaction of ionic bonds

▷ many weak interactions together = strong, hence importance & stability
 transient / statistical behaviors
 can lead to specific structures in self-assembled peptides (hydrophobicity)

□ Forces in mechanics



$E_{th} \sim 4 \cdot 10^{-21} \text{ J} \equiv \text{force} \cdot \text{distance}$
 over typical distance of nanometer (10^{-9} m)
 hence typical thermal force = $4 \cdot 10^{-12} \text{ N}$
 = 4 pN

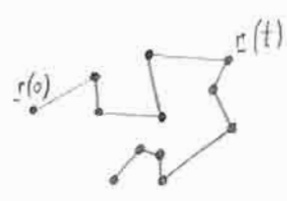
• measurements

thermal force	pN
optical tweezer	0.1 - 100 pN
magnetic tweezer	0.01 - 100 pN
AFM	10 - 10,000 pN

□ Case of thermal forces - modeling of DNA in solution - effect of kT

Thermal forces & diffusion

- Brownian motion (Brown 1827) in "simple" fluids (water)
 Jean Perrin quantified these observations (1900's)



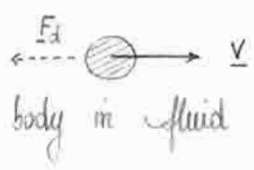
$\langle r \rangle = 0$ (ensemble average)
 $\langle \underline{r} \cdot \underline{r} \rangle$ linear in time = $4 D t$ ← diffusion coefficient
 = $\langle x^2 + y^2 \rangle$ in 2 dimensions

- what comes into D ? "guess" using dimensional analysis:

Fourier 1780: a physical equation must be { dimensionally consistent
 consistent in order (scalar, vector)
 e.g. $A = BC$
 $\downarrow \quad \downarrow \downarrow$
 $L/T \quad L T^{-1}$

{ dimensions: tell us about what we are meaning (length, velocity...)
 { units: a means to measure dimension (meters, meters per second...)

Postulate that D related to { the drag coefficient ξ
 the thermal energy $E_{th} (k_B T)$



drag force $\underline{F}_d = - \xi \underline{v}$
 $\xi = 6\pi \mu a$ (for sphere of radius a
 fluid of viscosity μ)

dimensions :

$$\left. \begin{array}{l} D : \frac{L^2}{t} \\ \Sigma : \frac{M}{t} \\ E_{th} : \frac{ML^2}{t} \end{array} \right\}$$

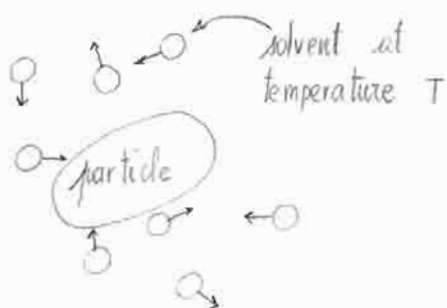
The dimensionally correct coefficient is

$$D \sim \frac{k_B T}{\Sigma}$$

never get coefficients from this approach

$$\langle x^2 \rangle \sim Dt$$

Brownian forces & the Langevin equation



model

solvent collisions on particle described by thermal force \underline{f}



Langevin equation : $\underline{F} = m \underline{a}$, $\underline{x}(t) =$ position

$$m \frac{d^2 \underline{r}(t)}{dt^2} = - \underbrace{\Sigma}_{\text{drag}} \frac{d\underline{r}(t)}{dt} + \underbrace{\underline{f}(t)}_{\text{thermal}} + \underbrace{\underline{G}(t)}_{\text{others: electric field, magnetic field}}$$

Thermal (Brownian) forces - fluctuate rapidly ($\sim 10^{-13}$ s in water)
- cannot be represented by a simple functional form

→ use stochastic (statistical) model for \underline{f}

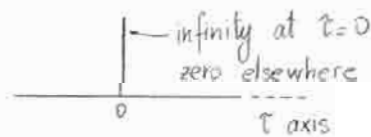
random in direction $\langle \underline{f} \rangle = 0$

uncorrelated on the time scale of particle motion

$$\langle \underline{f}(t) \cdot \underline{r}(t) \rangle = \langle \underline{f}(t) \rangle \langle \underline{r}(t) \rangle = 0$$

$$\left\{ \begin{array}{l} \langle \underline{f} \rangle = 0 \\ \langle \underline{f}(t) \cdot \underline{f}(t+\tau) \rangle = \underline{F} \delta(\tau) \end{array} \right.$$

$\delta(\tau)$ → delta function
 \underline{F} → unknown tensor



$$\int_{-\infty}^{+\infty} \delta(\tau) d\tau = 1$$

$$\int_{-\infty}^{+\infty} \delta(t) q(t) dt = q(0)$$

Manipulate equations with $\underline{\dot{r}} = \frac{d\underline{r}}{dt}$

$$* m \frac{d\underline{\dot{r}}}{dt} - \Sigma \underline{\dot{r}} + \underline{f} = 0$$

multiply by \underline{r} and take ensemble average.

$$m \left(\frac{d}{dt} \langle \underline{r} \underline{\dot{r}} \rangle - \langle \underline{\dot{r}} \underline{\dot{r}} \rangle \right) = - \Sigma \langle \underline{r} \underline{\dot{r}} \rangle + \underbrace{\langle \underline{r}(t) \underline{f}(t) \rangle}_{=0}$$

* equipartition theorem

$$\frac{1}{2} m \langle \underline{r} \dot{\underline{r}} \rangle = \frac{1}{2} k_B T \underline{\underline{\delta}} \quad \text{where} \quad \underline{\underline{\delta}} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$* m \frac{d}{dt} \langle \underline{r} \dot{\underline{r}} \rangle + \underline{\underline{\zeta}} \langle \underline{r} \dot{\underline{r}} \rangle = k_B T \underline{\underline{\delta}}$$

solve and use boundary conditions $\underline{r}(0) = 0$

$$\frac{d}{dt} \langle \underline{r} \dot{\underline{r}} \rangle = \frac{2k_B T}{\underline{\underline{\zeta}}} \underline{\underline{\delta}} \left(1 - \exp\left(-\frac{\underline{\underline{\zeta}}}{m} t\right) \right)$$

integrate and use BC $\dot{\underline{r}}(0) = 0$

$$\langle \underline{r} \dot{\underline{r}} \rangle = \frac{2k_B T}{\underline{\underline{\zeta}}} \underline{\underline{\delta}} \left\{ t + \frac{m}{\underline{\underline{\zeta}}} \left[\exp\left(-\frac{\underline{\underline{\zeta}}}{m} t\right) - 1 \right] \right\}$$

▷ 2 limiting cases

- $t \ll \frac{m}{\underline{\underline{\zeta}}}$ very short time scale

$$\exp\left(-\frac{\underline{\underline{\zeta}}}{m} t\right) \approx 1 - \frac{\underline{\underline{\zeta}}}{m} t + \frac{1}{2} \left(\frac{\underline{\underline{\zeta}}}{m}\right)^2 t^2 \dots$$

$$\langle \underline{r} \dot{\underline{r}} \rangle = \frac{k_B T}{m} \underline{\underline{\delta}} t^2 \quad \text{ballistic-like motion}$$

- $t \gg \frac{m}{\underline{\underline{\zeta}}}$ very long time scale

$$\langle \underline{r} \dot{\underline{r}} \rangle = \frac{2k_B T}{\underline{\underline{\zeta}}} \underline{\underline{\delta}} t \quad \text{diffusive motion}$$

▷ Some numbers:

- globular protein: }
 molecular weight 100 kDa
 $m = 170 \cdot 10^{-24} \text{ kg}$
 $\underline{\underline{\zeta}} \approx 6\pi \mu a$
 $a \approx 3 \text{ nm}$
 $\mu_{\text{water}} \approx 10^{-3} \text{ Pa}\cdot\text{s}$

- characteristic time $\frac{m}{\underline{\underline{\zeta}}} \approx 3 \cdot 10^{-12} \text{ s}$

- how far does it travel in time $m/\underline{\underline{\zeta}}$?

distance $\approx 10^{-11} \text{ m} = 10^{-2} \text{ nm}$

then diffusion Perrin model